

M597K: Solution to Homework 4

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1. Given a scalar $\Phi(x_1, x_2, x_3)$, does the gradient $\nabla\Phi = (\partial_{x_1}\Phi, \partial_{x_2}\Phi, \partial_{x_3}\Phi)$ satisfy the law of coordinate transformation for first-order tensors? That is,

$$\partial_{x'_i}\Phi'(x'_1, x'_2, x'_3) = \alpha_{i'k}\partial_{x_k}\Phi(x_1, x_2, x_3)?$$

Here $\alpha_{i'k}$ is the coordinate transformation from one rectangular coordinate system K to another rectangular coordinate system K' . Show your work. But you can skip this homework if you know how to solve the next problem.

Solution. (10 points each) Yes, it does. We have the inverse transformation

$$x_k = \alpha_{j'k}x'_j + x'_{0k}.$$

And the chain rule:

$$\frac{\partial}{\partial x'_i} = \frac{\partial x_k}{\partial x'_i} \frac{\partial}{\partial x_k} = \alpha_{i'k} \frac{\partial}{\partial x_k}.$$

Thus from $\Phi'(x'_1, x'_2, x'_3) = \Phi(x_1, x_2, x_3)$ and differentiation, we have

$$\begin{aligned} \partial_{x'_i}\Phi'(x'_1, x'_2, x'_3) &= \partial_{x'_i}\Phi(x_1, x_2, x_3) \\ &= \alpha_{i'k} \frac{\partial}{\partial x_k}\Phi(x_1, x_2, x_3). \end{aligned} \tag{1}$$

That is the law for transformation of first-order tensors.

2. Given a scalar function $\Phi = \Phi(x_1, x_2, x_3)$, do the quantities

$$\frac{\partial^2 \Phi}{\partial x_i \partial x_k}$$

form a tensor? Show your work.

Solution. Yes, it does. Continuing our solution in problem 1, we have

$$\begin{aligned} \partial_{x'_j} \partial_{x'_i} \Phi'(x'_1, x'_2, x'_3) &= \partial_{x'_j} \partial_{x'_i} \Phi(x_1, x_2, x_3) \\ &= \partial_{x'_j} \left[\alpha_{i'k} \frac{\partial}{\partial x_k} \Phi(x_1, x_2, x_3) \right] \\ &= \alpha_{j'm} \frac{\partial}{\partial x_m} \left[\alpha_{i'k} \frac{\partial}{\partial x_k} \Phi(x_1, x_2, x_3) \right] \\ &= \alpha_{j'm} \alpha_{i'k} \frac{\partial^2 \Phi(x_1, x_2, x_3)}{\partial x_m \partial x_k} \end{aligned} \tag{2}$$

which is the law for transformation of second-order tensors.

3. The stress tensor at a point has components given by

$$(p_{ij}) = \begin{pmatrix} 1 & -2 & 2 \\ -2 & 3 & 0 \\ 2 & 0 & -1 \end{pmatrix}.$$

Find the stress vector (\mathbf{p}_b) across an area normal to the unit vector

$$\mathbf{b} = (\mathbf{i}_1 - \mathbf{i}_2 + \mathbf{i}_3)/\sqrt{3}.$$

What is the normal stress across such an area (i.e, the projection $(\mathbf{p}_b \cdot \mathbf{b})\mathbf{b}$ of the vector \mathbf{p}_b on to \mathbf{b})?

Solution. A formula for \mathbf{p}_b is $\mathbf{p}_b = \mathbf{p}_i b_i$ where \mathbf{p}_i are the rows (or columns) of the matrix of the stress tensor. Thus

$$\mathbf{p}_b = (1, -2, 2)/\sqrt{3} + (-2, 3, 0)(-1)/\sqrt{3} + (2, 0, -1)/\sqrt{3} = (5, -5, 1)/\sqrt{3}.$$

The normal stress across such an area is

$$(\mathbf{p}_b \cdot \mathbf{b})\mathbf{b} = (11/3)\mathbf{b} = \frac{11\sqrt{3}}{9}(1, -1, 1).$$

4. For the stress tensor given in the previous problem,

(a) What is the total force on a unit disk whose normal is in the positive x_2 direction? *Answer:* the second column (or row) of the matrix (p_{ij}) .

(b) What is the x_3 component of the total force on a unit disk whose normal is in the positive x_1 direction? *Answer:* Either p_{13} or p_{31} , which are equal.

5. The unit base vectors \mathbf{i}'_i of a new coordinate system K' are given by

$$\mathbf{i}'_1 = \frac{\mathbf{i}_2 + \mathbf{i}_3}{\sqrt{2}}, \quad \mathbf{i}'_2 = \frac{\mathbf{i}_1 - \mathbf{i}_2 + \mathbf{i}_3}{\sqrt{3}}, \quad \mathbf{i}'_3 = \frac{2\mathbf{i}_1 + \mathbf{i}_2 - \mathbf{i}_3}{\sqrt{6}}.$$

The stress tensor p_{ik} in the system K is of the form

$$(p_{ik}) = \begin{pmatrix} p_1 & 0 & p_2 \\ 0 & 0 & 0 \\ p_2 & 0 & p_3 \end{pmatrix}.$$

Find the component p'_{13} of the stress tensor p'_{lm} in K' .

Solution. The $\alpha_{i'j} = \mathbf{i}'_i \cdot \mathbf{i}_j$ are found to be given by

$$(\alpha_{i'j}) = \begin{pmatrix} 0 & 1/\sqrt{2} & 1/\sqrt{2} \\ 1/\sqrt{3} & -1/\sqrt{3} & 1/\sqrt{3} \\ 2/\sqrt{6} & 1/\sqrt{6} & -1/\sqrt{6} \end{pmatrix}.$$

By formula of transformation, we have

$$\begin{aligned} p'_{13} &= \alpha_{1'k} \alpha_{3'm} p_{km} \\ &= \alpha_{1'1}(\alpha_{3'1}p_{11} + \alpha_{3'2}p_{12} + \alpha_{3'3}p_{13}) + \\ &\quad \alpha_{1'2}(\alpha_{3'1}p_{21} + \alpha_{3'2}p_{22} + \alpha_{3'3}p_{23}) + \\ &\quad \alpha_{1'3}(\alpha_{3'1}p_{31} + \alpha_{3'2}p_{32} + \alpha_{3'3}p_{33}) \\ &= 0(\dots) + \\ &\quad 1/\sqrt{2}(\alpha_{3'1}0 + \alpha_{3'2}0 + \alpha_{3'3}0) + \\ &\quad 1/\sqrt{2}(\alpha_{3'1}p_2 + \alpha_{3'2}0 + \alpha_{3'3}p_3) \\ &= 1/\sqrt{2}(\alpha_{3'1}p_2 + \alpha_{3'3}p_3) \\ &= 1/(2\sqrt{3})(2p_2 - p_3) = (2p_2 - p_3)/(2\sqrt{3}). \end{aligned} \tag{3}$$

Using the matrix multiplication $(p'_{ij}) = (\alpha)(p)(\alpha)^T$, one can find all components together. Here $(\alpha)^T$ denotes the transpose of (α) .

6. Let a_i , b_j , and c_k be the components of three vectors. Verify that the 27 quantities

$$d_{ijk} = 2a_i b_j c_k$$

form a tensor of order 3.

Solution. We know $a'_i = \alpha_{i'l} a_l$, $b'_j = \alpha_{j'm} b_m$, $c'_k = \alpha_{k'n} c_n$, thus

$$d'_{ijk} = 2a'_i b'_j c'_k = 2\alpha_{i'l} a_l \alpha_{j'm} b_m \alpha_{k'n} c_n = \alpha_{i'l} \alpha_{j'm} \alpha_{k'n} (2a_l b_m c_n) = \alpha_{i'l} \alpha_{j'm} \alpha_{k'n} d_{lmn}.$$

7. Form a scalar by contracting the tensor with the matrix

$$\begin{pmatrix} -5 & 0 & 1 \\ -1 & 3 & 7 \\ 4 & 8 & 2 \end{pmatrix}. \quad \text{Answer : } -5 + 3 + 2 = 0.$$

8. Given that

$$(T_{ik}) = \begin{pmatrix} 2 & -1 & 2 \\ 0 & 1 & 3 \\ -4 & 0 & 2 \end{pmatrix}, \quad \mathbf{A} = \mathbf{i}_1 - \mathbf{i}_2 + \mathbf{i}_3, \quad \mathbf{B} = 2\mathbf{i}_1 + \mathbf{i}_2 - \mathbf{i}_3.$$

Find the inner products $T_{ik}A_i$, $T_{ik}A_k$, and $T_{ik}A_iB_k$. If it is too tedious for you, you may choose i or k to be 2 if that i or k is not a dummy variable. (Note: First index in T is the row number.)

Solution. $T_{ik}A_i = T_{1k}A_1 + T_{2k}A_2 + T_{3k}A_3$.

For $k = 1$, we have $T_{i1}A_i = 2(1) + 0(-1) + (-4)(1) = -2$.

For $k = 2$, we have $T_{i2}A_i = (-1)1 + 1(-1) + 0(1) = -2$.

For $k = 3$, we have $T_{i3}A_i = 2(1) + 3(-1) + 2(1) = 1$.

For $i = 1$, $T_{ik}A_k = (2, -1, 2) \cdot (1, -1, 1) = 5$.

For $i = 2$, $T_{ik}A_k = (0, 1, 3) \cdot (1, -1, 1) = 2$.

For $i = 3$, $T_{ik}A_k = (-4, 0, 2) \cdot (1, -1, 1) = -2$.

You can use matrix multiplication to simplify the writing.

9. Show that the delta function $\delta_{ij} = \mathbf{i}_i \cdot \mathbf{i}_j$ satisfies the law of transformation for second-order tensors, where $\mathbf{i}_1, \mathbf{i}_2, \mathbf{i}_3$ are the unit vectors of a rectangular coordinate system K . This delta tensor is called the *unit tensor*.

Proof. Recall $\mathbf{i}'_k = \alpha_{k'm} \mathbf{i}_m$. Thus

$$\delta'_{ij} = \mathbf{i}'_i \cdot \mathbf{i}'_j = (\alpha_{i'm} \mathbf{i}_m) \cdot (\alpha_{j'n} \mathbf{i}_n) = \alpha_{i'm} \alpha_{j'n} \mathbf{i}_m \cdot \mathbf{i}_n = \alpha_{i'm} \alpha_{j'n} \delta_{mn}.$$

So the delta satisfies the law of transformation for second-order tensors.

10. (Optional) In nonlinear elasticity, the strain tensor is more accurately defined to be

$$u_{ik} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_k} + \frac{\partial u_k}{\partial x_i} + \frac{\partial u_l}{\partial x_i} \frac{\partial u_l}{\partial x_k} \right)$$

where u_i are again components of the displacement vector. Show it still satisfies the transformation law for second-order tensors.

(This set of homework looks a lot, but try your best. Some of them are really easy.)

Proof. Recall $u'_i = \alpha_{i'k} u_k$, and

$$u'_{ik} = \frac{1}{2} \left(\frac{\partial u'_i}{\partial x'_k} + \frac{\partial u'_k}{\partial x'_i} + \frac{\partial u'_l}{\partial x'_i} \frac{\partial u'_l}{\partial x'_k} \right),$$

and

$$\frac{\partial}{\partial x'_i} = \frac{\partial x_j}{\partial x'_i} \frac{\partial}{\partial x_j} = \alpha_{i'j} \frac{\partial}{\partial x_j}.$$

Thus

$$\begin{aligned} 2u'_{ik} &= \frac{\partial u'_i}{\partial x'_k} + \frac{\partial u'_k}{\partial x'_i} + \frac{\partial u'_l}{\partial x'_i} \frac{\partial u'_l}{\partial x'_k} \\ &= \alpha_{k'm} \frac{\partial}{\partial x_m} (\alpha_{i'n} u_n) + \alpha_{i'n} \frac{\partial}{\partial x_n} (\alpha_{k'm} u_m) + \\ &\quad + \alpha_{i'n} \frac{\partial}{\partial x_n} (\alpha_{l's} u_s) \alpha_{k'm} \frac{\partial}{\partial x_m} (\alpha_{l'r} u_r) \\ &= \alpha_{k'm} \alpha_{i'n} \frac{\partial u_n}{\partial x_m} + \alpha_{i'n} \alpha_{k'm} \frac{\partial u_m}{\partial x_n} + \\ &\quad \alpha_{i'n} \alpha_{k'm} \alpha_{l's} \alpha_{l'r} \frac{\partial u_s}{\partial x_n} \frac{\partial u_r}{\partial x_m} \\ &= \alpha_{k'm} \alpha_{i'n} \left(\frac{\partial u_n}{\partial x_m} + \frac{\partial u_m}{\partial x_n} + \delta_{rs} \frac{\partial u_s}{\partial x_n} \frac{\partial u_r}{\partial x_m} \right) \\ &= \alpha_{k'm} \alpha_{i'n} \left(\frac{\partial u_n}{\partial x_m} + \frac{\partial u_m}{\partial x_n} + \frac{\partial u_r}{\partial x_n} \frac{\partial u_r}{\partial x_m} \right) \\ &= \alpha_{i'n} \alpha_{k'm} (2u_{nm}), \end{aligned} \tag{4}$$

where we've used the orthogonality condition $\alpha_{l's} \alpha_{l'r} = \delta_{rs}$. That is the transformation law for second-order tensors.